

# SHRINKING SOLIDS

COCURRENT and PLUG FLOW  
COUNTERCURRENT

Shrinking solids



Plug Flow

$W(g)$  = mass of solids in the bed

$\rho_0 (g/cm^3)$  = density of solids

$C_{B0} (mol/cm^3)$  = molar density of solids

$X_B = \frac{R_0^3 - R^3}{R_0^3}$  = fractional conversion

$F_{B0} (\frac{mol B}{s})$  = molar feed rate of solids

$R_B (\frac{mol B}{cm^3 \text{ solids } s})$  = reaction rate of B

$$F_{B0} \frac{dX_B}{dW} = \frac{R_B}{\rho_0}$$

$$R_B = \frac{4\pi R^2 C_{B0} \left(-\frac{dR}{dt}\right)}{\frac{4}{3}\pi R^3} = \frac{3 C_{B0} R}{R \rho_0}$$

$$\left( \frac{cm}{s} \right) \left. \begin{aligned} -R &= + \frac{dR}{dt} = -k C_{A0}^m R^2 \\ \end{aligned} \right\}$$

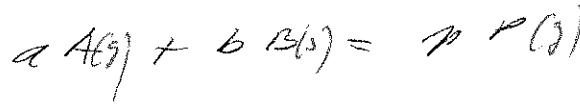
$d \geq 0$  kinetic control

$d \leq \frac{1}{2}$  diffusion control  
- Fickian  
- Cussler

$d = 1$

shrink

(New 1991)



Plug Flow Solids — Stirring vessels

Plug Flow Gas (liquid) — Cocurrent

spherical

Rate law

$$\frac{dR}{dt} = -k C_A^n R^2 = -R$$

$$R = -\frac{dR}{dt}$$

$$R > 0$$

$$F_{B0} \frac{dX_B}{dW} = \frac{3C_{A0} R}{R_0}$$

$$X_B = 1 - \left(\frac{R}{R_0}\right)^3$$

$$dX_B = -\frac{3R^2}{R_0^3} dR$$

$$\frac{3F_{B0} R^2}{R_0^3} \frac{dR}{dW} = -\frac{3C_{A0} R}{R_0}$$

$$\frac{dR}{dW} = -\frac{R_0^3 C_{A0} R}{F_{B0} R^3} ; W=0 \quad R=R_0$$

$$-\frac{d}{dW} (Q_g C_A) - \frac{3aC_{A0} R}{6R_0} = 0$$

$$\frac{d}{dW} (Q_g C_A) = -\frac{3}{R} \left(\frac{aC_{A0}}{6R_0}\right) R ; W=0 \quad Q_g C_A = Q_{g0} C_{A0}$$

$$\frac{d}{dW} \left[ F_{B0} X_B + \overset{F_A}{Q_g C_A} \right] = 0$$

$$Q_g C_A + \frac{a}{b} F_{B0} X_B = Q_{g0} C_{A0}$$

$$C_A = \frac{Q_{g0}}{Q_g} C_{A0} - \frac{a}{b} \frac{F_{B0}}{Q_g} X_B$$

$$C_A = \frac{Q_{g0}}{Q_g} C_{A0} - \frac{a F_{B0}}{b (Q_g)'} X_B$$

$$C_A = C_{A0} \frac{Q_{g0}}{Q_g} \left[ 1 - \frac{a F_{B0}}{b F_{A0}} X_B \right]$$

$$C_A = C_{A0} \left( \frac{Q_{g0}}{Q_g} \right) [1 - \beta X_B]$$

$$\beta = \frac{a F_{A0}}{b F_{A0}}$$

Ideal gas law

$$\frac{Q_g}{Q_{g0}} = \frac{T}{T_0} \frac{P_0}{P} (1 + \epsilon_A X_A) = \frac{T P_0}{T_0 P} (1 + \beta X_B)$$

$$X_A = \frac{F_{A0} - F_A}{F_{A0}}$$

$$\epsilon_A = \frac{(\sum \nu_i^0)}{(-\nu_A)} y_{A0}$$

$$\frac{d}{dW} \left[ F_{B0} X_B \frac{a}{b} + F_A \right] = 0$$

$$F_A + F_{B0} X_B \frac{a}{b} = F_{A0}$$

$$F_{A0} - F_A = F_{A0} X_A = F_{B0} X_B \frac{a}{b}$$

$$X_A = \frac{a F_{B0}}{b F_{A0}} X_B = \beta X_B$$

$$C_A = C_{A0} \frac{T_0 P}{T P_0} \frac{1 - \beta X_B}{1 + \epsilon_A \beta X_B}$$

$$X_B = 1 - \left( \frac{P}{P_0} \right)^3$$

Assume  $\epsilon_A = 0$   
 $T = \text{const}$   
 $P = \text{const}$

$$C_A = C_{A0} (1 - \beta X_B)$$

$\epsilon_A = 0$  converted to

$$\frac{dR}{dW} = - \frac{R_0^3 C_{A0}}{F_{B0} \beta} \frac{R}{R^3}$$

gm  
phero

$$\frac{dR}{dW} = - \left( \frac{R_0^3 C_{A0}}{F_{B0} \beta} \right) \frac{k C_A^m R^2}{R^3}$$

$$\frac{dR}{dW} = - \left( \frac{R_0^3 C_{A0} k}{F_{B0} \beta} \right) C_{A0}^m (1 - \beta X_B)^m R^{\alpha-3}$$

$$\frac{dR}{dW} = - \left( \frac{R_0^3 C_{A0} k C_{A0}^m}{F_{B0} \beta} \right) \left[ 1 - \beta + \beta \left( \frac{R}{R_0} \right)^3 \right] R^{\alpha-3}$$

$$W = \frac{F_{B0} \beta}{C_{A0} k C_{A0}^m R_0^3} \frac{1}{R_0^3} \int_R^{R_0} \frac{dR}{\left[ 1 - \beta + \beta \left( \frac{R}{R_0} \right)^3 \right] R^{\alpha-3}}$$

$$\frac{F_{B0} \beta}{C_{A0}} = \dot{m}_{B0} \left( \frac{g}{s} \right)$$

$\epsilon_{\text{space}}$   
 $\frac{1}{\text{space}}$

$$\frac{W}{\dot{m}_{B0}} = \frac{1}{k C_{A0}^m R_0^3} \int_R^{R_0} \frac{dR}{\left[ 1 - \beta + \beta \left( \frac{R}{R_0} \right)^3 \right] R^{\alpha-3}}$$

- Large external gas  $C_{A0} = \text{const}$   $\beta \Rightarrow 0$
- Stochastic metric  $\alpha = 1$   $\beta = 1$

Large excess of gas:  $\beta \rightarrow 0$

$$\frac{W}{\dot{m}_{B0}} = \frac{1}{k_{CA0}^n R_0^3} \int_R^{R_0} R^{3-\alpha} dR$$

$$= \frac{R_0^{3-\alpha+1} - R^{3-\alpha+1}}{k_{CA0}^n R_0^3 (3-\alpha+1)} = \frac{R_0^{4-\alpha} - R^{4-\alpha}}{k_{CA0}^n R_0^3 (4-\alpha)}$$

$$\frac{W}{\dot{m}_{B0}} = \frac{R_0^{4-\alpha}}{k_{CA0}^n (4-\alpha)} \left[ 1 - \left( \frac{R}{R_0} \right)^{4-\alpha} \right]$$

Stoichiometric ratio  $\beta \Rightarrow 1$

$$\frac{W}{\dot{m}_{B0}} = \frac{1}{k_{CA0}^n} \int_R^{R_0} R^{-\alpha} dR = \frac{R_0^{1-\alpha} - R^{1-\alpha}}{k_{CA0}^n (1-\alpha)}$$

$$\frac{W}{\dot{m}_{B0}} = \frac{R_0^{1-\alpha}}{k_{CA0}^n (1-\alpha)} \left[ 1 - \left( \frac{R}{R_0} \right)^{1-\alpha} \right]$$

$$\frac{W_{\beta=0}}{W_{\beta=1}} = \frac{1-\alpha}{4-\alpha} \frac{\left[ 1 - \left( \frac{R}{R_0} \right)^{4-\alpha} \right]}{\left[ 1 - \left( \frac{R}{R_0} \right)^{1-\alpha} \right]}$$

for complete reaction  $R=0$

$$\frac{W_{\beta=0}}{W_{\beta=1}} = \frac{1-\alpha}{4-\alpha} = \begin{cases} \frac{1}{4} & \alpha=0 \\ \frac{2}{5} & \alpha=-1 \\ \frac{1}{2} & \alpha=-2 \end{cases}$$

$$\frac{dB}{dt} = -kA^\alpha R^\alpha$$


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$$\alpha = 0$$

$$\tau_{B=0} = \frac{R_0}{4kC_{A0}^m}$$

$$\tau_{B=1} = \frac{R_0}{kC_{A0}^m}$$

$$\frac{\tau_{B=1}}{\tau_{B=0}} = 4$$


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$$\alpha = -\frac{1}{2}$$

$$\tau_{B=0} = \frac{R_0^{3/2}}{\frac{9}{2}kC_{A0}} = \frac{2R_0^{3/2}}{9kC_{A0}}$$

$$\tau_{B=1} = \frac{R_0^{3/2}}{\frac{3}{2}kC_{A0}} = \frac{2R_0^{3/2}}{3kC_{A0}}$$

$$\frac{\tau_{B=1}}{\tau_{B=0}} = 3$$


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$$\alpha = -1$$

$$\tau_{B=0} = \frac{R_0^2}{5kC_{A0}^m}$$

$$\tau_{B=1} = \frac{R_0^2}{2kC_{A0}^m}$$

$$\frac{\tau_{B=1}}{\tau_{B=0}} = \frac{5}{2} = 2.5$$


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$$\alpha = 1$$

$$\tau_{B=0} = \frac{1}{3kC_{A0}^m}$$

$$\tau_{B=1} = \infty$$

$$\frac{\tau_{B=1}}{\tau_{B=0}} = \infty$$

$$\alpha = \frac{1}{2} \quad \tau_{B=0} = \frac{R_0^{1/2}}{kC_{A0}^{3/2}} = \frac{2}{3} \quad \tau_{B=1} = \frac{R_0^{1/2}}{\frac{1}{2}kC_{A0}^{3/2}} = 2 \quad \frac{\tau_{B=1}}{\tau_{B=0}} = 3$$

$\alpha = 0$  *isotropic* *case*

To get  $R = 0$

$\beta \Rightarrow 0$

$$\tau_{W, \beta=0} = \frac{W}{\dot{m}_{\beta=0}} = \frac{R_0}{4 \sqrt{C_{D0}^*}}$$

*in oxygen gas*

$\beta \Rightarrow 1$

$$\tau_{W, \beta=1} = \frac{W}{\dot{m}_{\beta=1}} = \frac{R_0}{4 C_{D0}^*}$$

*2nd*

$$\frac{\tau_{W, \beta=0}}{\tau_{W, \beta=1}} = \frac{1}{4} = 0.25$$

*1/4 size of bed required*

$$\alpha = -\frac{1}{2} \quad \tau_{W, \beta=0} = \frac{R_0^{3/2}}{9 \sqrt{C_{D0}^*}} = \frac{2 R_0^{3/2}}{9 \sqrt{C_{D0}^*}} \quad \frac{\tau_{W, \beta=0}}{\tau_{W, \beta=1}} = \frac{2}{9} = 0.222$$

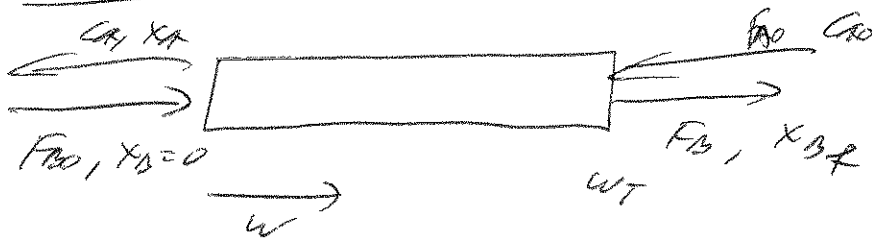
$$\tau_{W, \beta=1} = \frac{R_0^{3/2}}{3 \sqrt{C_{D0}^*}} = \frac{2 R_0^{3/2}}{3 \sqrt{C_{D0}^*}}$$

$\alpha = -1$  *different* *case*,  $R = 0$

$$\tau_{W, \beta=0} = \frac{R_0^2}{5 \sqrt{C_{D0}^*}}$$

$$\tau_{W, \beta=1} = \frac{R_0^2}{2 \sqrt{C_{D0}^*}}$$

$$\frac{\tau_{W, \beta=0}}{\tau_{W, \beta=1}} = \frac{2}{5} = 0.4$$

COUNTER CURRENT FLOW

$$F_{B0} \frac{dX_B}{dW} = \frac{3 C_{B0}}{R P_0} R$$

$$W = 0 \quad X_B = 0$$

$$+ \frac{d}{dW} (Q_g C_A) = + \frac{dF_A}{dW} = \frac{3 a C_{A0}}{R b C_B} R$$

$$W = W_T \quad C_A = C_{A0}, \quad F_A = F_{A0}$$

$$\frac{d}{dW} \left[ F_A - \frac{a}{b} F_{B0} X_B \right] = 0$$

$$F_A - \frac{a}{b} F_{B0} X_B = F_{A0} - \frac{a}{b} F_{B0} X_{BF}$$

$$F_A = F_{A0} - \frac{a}{b} F_{B0} (X_{BF} - X_A)$$

$$F_A = F_{A0} \left[ 1 - \beta (X_{BF} - X_A) \right]$$

$$\beta = \frac{a F_{B0}}{b F_{A0}}$$

$$C_A = C_{A0} \frac{Q_{S0}}{Q_g} \left[ 1 - \beta (X_{BF} - X_A) \right]$$

$$\frac{Q_g}{Q_{S0}} = \frac{T P_0}{T_0 P} (1 + E_A X_A)$$

$$X_A = \beta (X_{BF} - X_A)$$



$$C_A = C_{A0} \frac{T_0 P}{T P_0} \frac{1 - \beta (x_{A2} - x_A)}{1 + E_A \beta (x_{A2} - x_A)}$$

$$x_A = 1 - \left(\frac{R}{R_0}\right)^3$$

$$dx_A = -\frac{3R^2}{R_0^3} dR$$

Let  $T = \text{const}$   $P = \text{const}$   $E_A = 0$

$$C_A = C_{A0} [1 - \beta (x_{A2} - x_A)]$$

$$\frac{dR}{dW} = - \left( \frac{R_0^3 C_{A0}}{F_{A0} \beta} \right) \frac{k C_A^n R^{\alpha}}{R^3}$$

$$\frac{dR}{dW} = - \frac{R_0^3 C_{A0} k C_{A0}^n}{F_{A0} \beta} [1 - \beta (x_{A2} - x_A)] R^{\alpha-3}$$

$$x_{A2} = 1 - \left(\frac{R_2}{R_0}\right)^3 \quad x_A = 1 - \left(\frac{R}{R_0}\right)^3$$

$$x_{A2} - x_A = \left(\frac{R}{R_0}\right)^3 - \left(\frac{R_2}{R_0}\right)^3$$

$$\frac{dR}{dW} = - \left( \frac{R_0^3 C_{A0} k C_{A0}^n}{F_{A0} \beta} \right) \left[ 1 + \beta \left(\frac{R_2}{R_0}\right)^3 - \beta \left(\frac{R}{R_0}\right)^3 \right] R^{\alpha-3}$$

$$\frac{F_{A0} \beta}{C_{A0}} = \alpha R_0$$

$$W=0 \quad R=R_0$$

$$W=W_T \quad R=R_2$$

$$\frac{W}{\dot{m}_{B0}} = \frac{1}{kA_0^m R_0^3} \int_R^{R_0} \frac{dR}{\left[1 + \beta \left(\frac{R}{R_0}\right)^3 - \beta \left(\frac{R}{R_0}\right)^3\right] R^{\alpha-3}}$$

$$\frac{W_T}{\dot{m}_{B0}} = \frac{1}{kA_0^m R_0^3} \int_{R_f}^{R_0} \frac{dR}{\left[1 + \beta \left(\frac{R}{R_0}\right)^3 - \beta \left(\frac{R}{R_0}\right)^3\right] R^{\alpha-3}}$$

Large exothermic  $\beta \rightarrow 0$

$$\begin{aligned} \frac{W_T}{\dot{m}_{B0}} &= \frac{1}{kA_0^m R_0^3} \int_{R_f}^{R_0} R^{3-\alpha} dR \\ &= \frac{R_0^{4-\alpha}}{kA_0^m (4-\alpha)} \left[1 - \left(\frac{R_f}{R_0}\right)^{4-\alpha}\right] \quad Q50 \end{aligned}$$

same as counter

Stoichiometric rate  $\beta=1$

$$\frac{W_T}{\dot{m}_{B0}} = \frac{1}{kA_0^m R_0^3} \int_{R_f}^{R_0} \frac{dR}{\left[1 + \left(\frac{R}{R_0}\right)^3 - \left(\frac{R}{R_0}\right)^3\right] R^{\alpha-3}}$$

For complete reaction of particles  $R_f = 0$

$$\begin{aligned} \frac{W_{T, \text{comp}}}{\dot{m}_{B0}} &= \frac{1}{kA_0^m} \int_0^{R_0} \frac{dR}{(R_0^3 - R^3) R^{\alpha-3}} \\ &= \frac{1}{kA_0^m} \int_0^1 \frac{dR}{(R_0^3/R^3 - 1) R^{\alpha-2}} \end{aligned}$$

Compare cocurrent and counter current

$$\text{let } R = \frac{R}{R_0} \leq 1$$

$$\frac{dR}{R^2} = \frac{R_0 dR}{R_0^2 R^2} = R_0^{1-\alpha} \frac{dR}{R^2}$$

CO CURRENT FLOW (REACTION SIZE FOR COMPLETE CONVERSION)

$\beta = 1$  stock

$$\frac{W_{CCF \text{ comp}}}{m_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 R^{-2} dR = \frac{R_0^{1-\alpha}}{k C_{A0}^m (1-\alpha)}$$

COUNTER CURRENT FLOW

$$\frac{W_{CCF \text{ comp}}}{m_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 \frac{R^{-2} dR}{\left(\frac{1}{R^3} - 1\right)} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 \frac{R^{3-\alpha} dR}{1-R^3}$$

$$\text{But } \frac{1}{R} \geq 1 \quad \frac{1}{R} - 1 > 0$$

$$0 \leq R \leq 1$$

$$\frac{W_{CCF \text{ comp}}}{m_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 R^{3-\alpha} \sum_{n=0}^{\infty} \frac{(R^3)^n}{R^3} dR$$

$$= \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 \sum_{n=0}^{\infty} R^{3n+3-\alpha} dR$$

$$= \frac{R_0^{1-\alpha}}{k C_{A0}^m} \sum_{n=0}^{\infty} \left( \frac{1}{3n+4-\alpha} \right) \quad \text{not convergent?}$$

CCF-5

$$\frac{W_{CCF \text{ case}}}{\sin \alpha} = \frac{R_0^{1-\alpha}}{kL_0^{\alpha}} \left[ \frac{1}{4-\alpha} + \frac{1}{7-\alpha} + \frac{1}{10-\alpha} + \frac{1}{13-\alpha} + \frac{1}{16-\alpha} + \dots \right]$$

look again

$$\frac{W_{T \text{ case}}}{\sin \alpha} = \frac{R_0^{1-\alpha}}{kL_0^{\alpha}} \int_0^1 \frac{r^{3-\alpha}}{1-r^3} dr$$

~~let  $(1-r^3) = x$~~

$x = r^3$

$dx = 3r^2 dr$

$r = x^{1/3}$

$dr = \frac{1}{3} \frac{1}{x^{2/3}} dx$

$$\int_0^1 \frac{r^{3-\alpha}}{1-r^3} dr = \int_0^1 \frac{(x^{1/3})^{(3-\alpha)} dx}{3 x^{2/3} (1-x)} = \frac{1}{3} \int_0^1 \frac{x^{1-\frac{\alpha}{3}}}{x^{2/3} (1-x)} dx$$

$$= \frac{1}{3} \int_0^1 \frac{x^{1/3 - \frac{\alpha}{3}}}{1-x} dx$$

$$= \frac{1}{3} \int_0^1 \frac{x^{(1/3 - \alpha/3)}}{1-x} dx = \frac{1}{3} \int_0^1 \frac{x^{\frac{1-\alpha}{3}}}{1-x} dx = \frac{1}{3} \int_0^1 x^{\frac{1-\alpha}{3}} \sum x^n dx$$

$$= \frac{1}{3} \int_0^1 \sum x^{\frac{3n+1-\alpha}{3}} dx = \frac{1}{3} \sum \frac{1}{3n+4-\alpha}$$

The problem is that at  $w=0$   
 fresh feeds at  $R=R_0$ ,  $x_B=0$   
 meet total react gas at  $r=0$ , i.e.  $C_A=0$   
 This causes at the reactor end  
 a singularity for feeds

$$\frac{dR}{dt} \Big|_{w=0} = -0 \quad \text{or} \quad \frac{dt}{dR} = \infty!$$

If we do not quite react the  
 feeds completely then the process  
 is feasible, say  $\frac{R_f}{R_0} = \epsilon$  ( $\epsilon < 1$ ,  $\epsilon > 0$ )  
 then  $x_{Bf} = 1 - \epsilon^3$

$$\frac{W_T}{\dot{m}_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_{\epsilon}^1 \frac{R^{3-\alpha} dR}{(1 + \epsilon^3 - R^3)}$$

and there is no singularity at  $R=1$

Alternatively we can take  $bF_{A0} > aF_{B0}$   
 say slight excess  $\beta = 1 - \epsilon$

$$\frac{W_T}{\dot{m}_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_{\beta}^1 \frac{R^{3-\alpha} dR}{(1 + \beta R^3 - \beta R^3)}$$

Now we can take  $R_f = 0$  if  $\beta = 1 - \epsilon$

$$\frac{W_T}{\dot{m}_{B0}} = \frac{R_0^{1-\alpha}}{k C_{A0}^m} \int_0^1 \frac{R^{3-\alpha} dR}{1 - (1 - \epsilon) R^3}$$

and there is no singularity at  $R=1$